

# Diagrammatic Expansion for Positive Spectral Functions in the Steady-State Limit



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## Abstract

Recently, a method was presented [1] for constructing self-energies within many-body perturbation theory that are guaranteed to produce a positive spectral function for equilibrium systems, by representing the self-energy as a product of half-diagrams on the forward and backward branches of the Keldysh contour. We derive an alternative half-diagram representation that is based on products of retarded diagrams. Our approach extends the method to systems out of equilibrium. When a steady-state limit exists, we show that our approach yields a positive definite spectral function in the frequency domain.

## 1. The Problem

- Some approximations for the self-energy can lead to negative valued spectral functions.
- This prevents the probability interpretation of the spectral function, and can lead to instability in self-consistent calculations.
- This issue arises for example for:

$$\Sigma_{xc}(z_1, z_2) = \text{Diagram 1} + \text{Diagram 2} \quad (1)$$

## 2. Self-Energy in Terms of Half-Diagrams

- In Ref. [1] it was shown that in zero temperature the spectral function will be positive semi-definite (PSD) if the self-energy can be represented as a sum of squares of half-diagrams.
- For example

$$\text{Diagram 1} = \text{Diagram 2} \times \text{Diagram 3} \quad (2)$$

- This splitting is obtained by placing a complete basis set at the end-point of the contour at  $t = \infty$ :

$$\text{Diagram 1} = \sum_i |\chi_i\rangle\langle\chi_i| \quad (3)$$

- The half-diagrams are then (anti)-time ordered.
- This requires the assumption that the initial state  $|\Phi_0\rangle$  is connected to the state obtained in the distant future:

$$\mathcal{U}(\infty, -\infty)|\Phi_0\rangle = e^{i\alpha}|\Phi_0\rangle. \quad (4)$$

## 3. Generalization to Steady-State

- In Ref. [2] we show that the same proof can be performed without the assumption in Eq.(4), by placing the basis set at  $t = -\infty$ :

$$\text{Diagram 1} = \sum_i |\chi_i\rangle\langle\chi_i| \quad (5)$$

- The PSD nature of the spectral function can now be shown in the steady-state limit.
- In this case the half-diagrams are no longer time-ordered, but retarded. For example

$$\text{Diagram 1} = \text{Diagram 2} \times \text{Diagram 3} \quad (6)$$

where the circling of vertices denotes a retarded piece, the other vertices being retarded with respect to the vertex marked with a double circle.

- In a general half-diagram all internal vertices are retarded with respect to vertex 1 (or 2). This requires defining a general retarded diagram [2].

## 4. Example: Second Born Approximation

The 2B approximation is PSD, since

$$\text{Diagram 1} + \text{Diagram 2} = \text{Diagram 3} + \text{Diagram 4} = \frac{1}{2} \left[ \text{Diagram 5} + \text{Diagram 6} \right]^2 \quad (7)$$

## 5. Example: $gW$ Approximation

The  $gW$  approximation with the RPA polarization function  $P(1,2) = g(1,2)g(2,1)$  is PSD, since using

$$W^{\lessgtr}(1,2) = \int d3d4 W^R(1,3) P^{\lessgtr}(3,4) W^A(4,2) \quad (8)$$

we can write the xc-self-energy as

$$\Sigma_{xc}^{\lessgtr} = \text{Diagram 1} = \text{Diagram 2} \times \text{Diagram 3} \quad (9)$$

## 6. Generalized Retarded Compositions

- In order to define general retarded compositions, we will first define general contour ordered components
- A contour function that is representable as a diagram can be replaced by a real-time function whenever the order along the contour  $\gamma$  is fixed. For example

$$\bar{\Sigma}_{xc}(z_a, z_b, z_c, z_d) = \bar{\Sigma}_{xc}^{1234}(t_a, t_b, t_c, t_d) \text{ when } z_a > z_b > z_c > z_d. \quad (10)$$

- For a two-point function these components are equal to the greater and lesser components:

$$G^{12}(z_a, z_b) = G^>(z_a, z_b), \quad G^{21}(z_a, z_b) = G^<(z_a, z_b) \quad (11)$$

This set of functions for each permutation of the arguments encodes all the information in the original contour function.

### 6.1 One Vertex

Take for example a half-diagram with a single internal vertex (with diagrammatic representation on the right):

$$D(z_a) = \int_{\gamma} dz_b D(z_a, z_b) \quad \text{Diagram 1} = \text{Diagram 2} \quad (12)$$

If one defines a retarded composition

$$D^R(t_a, t_b) = \theta(t_a - t_b) (D^>(t_a, t_b) - D^<(t_a, t_b)) = \theta_{ab} O^{[1,2]}(t_a, t_b), \quad (13)$$

it follows that

$$D^1(t_a) = \int_{t_0}^{\infty} dt_b D^R(t_a, t_b) \quad \text{Diagram 1} = \text{Diagram 2} \quad (14)$$

### 6.2 Two Vertices

A half-diagram with two internal vertices

$$D(z_a) = \int_{\gamma} dz_b dz_c D(z_a, z_b, z_c) \quad \text{Diagram 1} = \text{Diagram 2} \quad (15)$$

motivates the definition of a retarded composition for which

$$D^1(t_a) = \int_{t_0}^{\infty} dt_b dt_c D^{R(1,23)}(t_a, t_b, t_c) \quad \text{Diagram 1} = \text{Diagram 2} \quad (16)$$

This is achieved for example by

$$D^{R(1,23)} = \theta_{abc} D^{[[1,2],3]} + \theta_{acb} O^{[[1,3],2]}. \quad (17)$$

### 6.3 N Vertices

For a general half-diagram with  $N$  internal vertices

$$D(z_{n_1}) = \int_{\gamma} dz_{n_2} \dots dz_{n_N} D(z_{n_1}, \dots, z_{n_N}) \quad \text{Diagram 1} = \text{Diagram 2} \quad (18)$$

one then has

$$D^1(t_{n_1}) = \int_{t_0}^{\infty} dt_{n_2} \dots dt_{n_N} D^{R(1,2 \dots n)}(t_{n_1}, \dots, t_{n_N}) \quad \text{Diagram 1} = \text{Diagram 2} \quad (19)$$

with the retarded composition

$$D^{R(1,2 \dots n)} = \sum_{P \in S_{n-1}} \theta_{1P(2) \dots P(n)} D^{[1, P(2), \dots, P(n)]} \quad (20)$$

## References

- [1] G. Stefanucci, Y. Pavlyukh, A.M. Uimonen, and R. van Leeuwen, Phys. Rev. B 90, 115134 (2014).  
 [2] M. J. Hyrkäs, D. Karlsson, and R. van Leeuwen *Contour calculus for many-particle functions*, to be published